

# Newtonian Noise

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# Overview

What is Newtonian noise?

Seismic sources

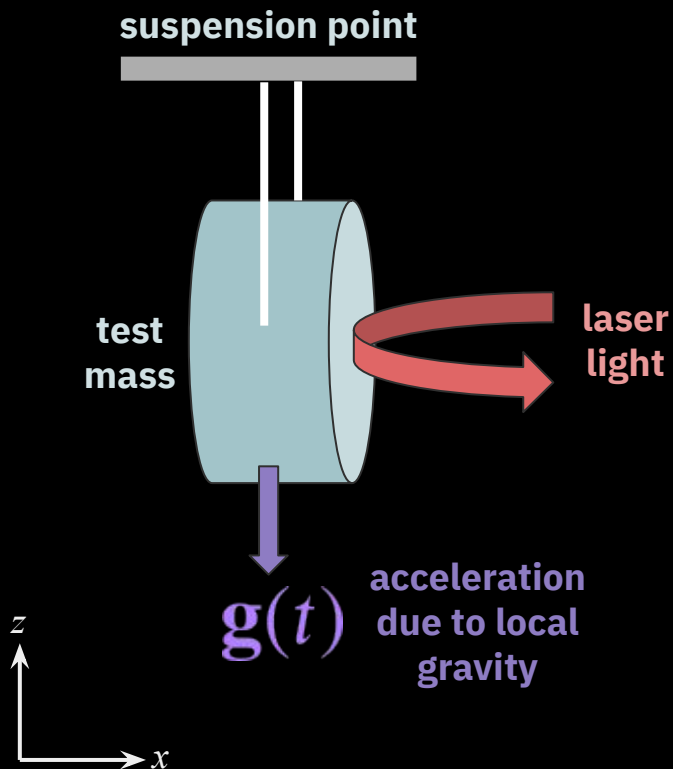
Atmospheric sources

Mitigation schemes

Implementation in pygwinc

# Newtonian noise

(a.k.a. gravity gradient noise)



$$\mathbf{g}(t) = \mathbf{g}_0 + \delta\mathbf{g}(t)$$

large, static  
component

fluctuating  
component  
(Newtonian  
noise)

$$\ddot{h}(t) = \frac{\delta g_x(t)}{L} \Leftrightarrow h(\omega) = \frac{\delta g_x(\omega)}{-\omega^2 L}$$

# Computing NN

Newton's law of universal gravitation:

$$\mathbf{g}(t) = G_N \int d^3\mathbf{r} \frac{\rho(\mathbf{r}, t)}{r^2} \hat{\mathbf{r}}$$

NN can be caused by

- the **motion of a point mass**,
- a **change in the density distribution in an extended mass**, or
- a **displacement of the boundary of an extended mass**, even in the absence of a change in density.

# What causes NN?

## Continuous sources:

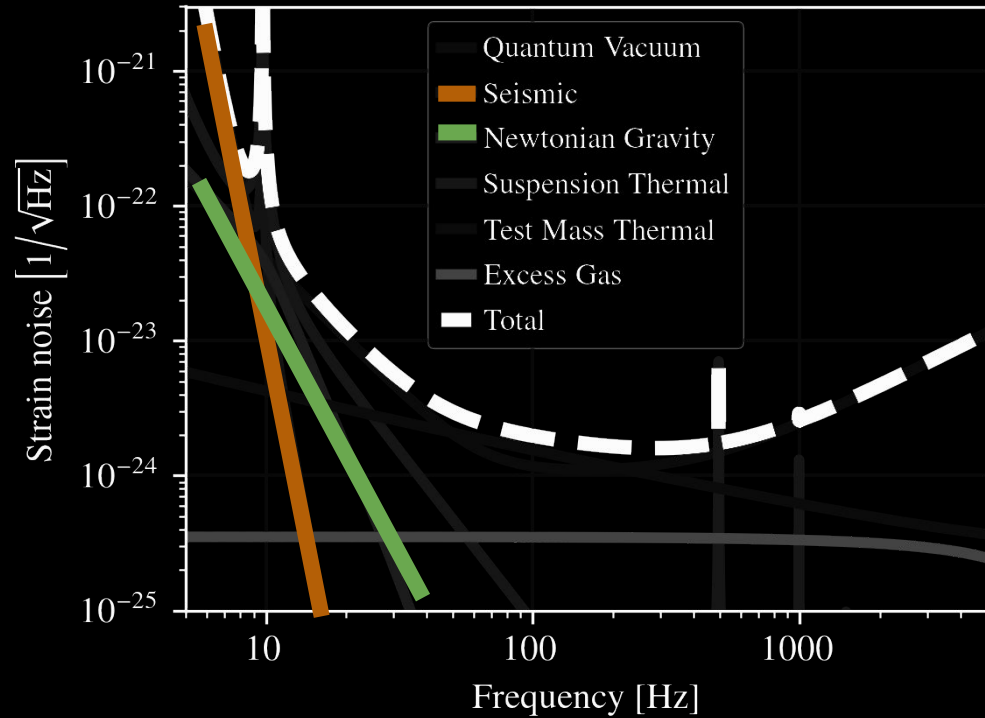
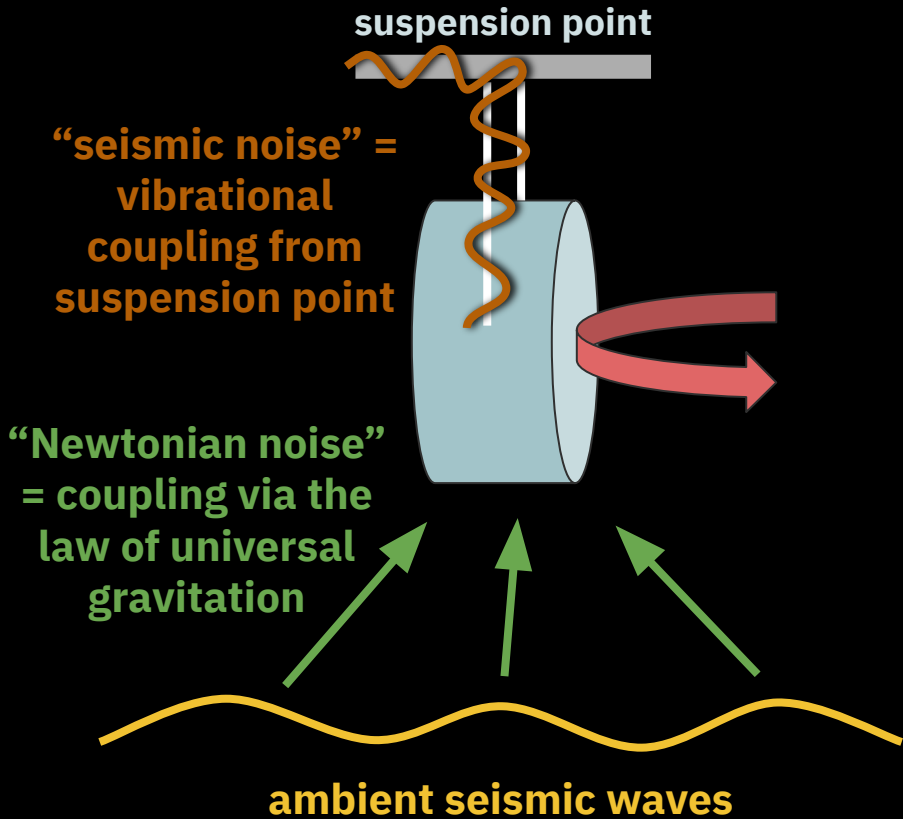
- Ambient seismic waves
  - Ambient acoustic waves
  - Other atmospheric processes
  - Flowing or bubbling liquids
- } implemented in pygwinc

## Transient sources:

- Earthquakes
- Shock waves
- Cars and airplanes
- Humans [Thorne & Winstein (1998) [PRD 60 082001](#)]
- Tumbleweeds [Creighton (2008) [COG 25\(12\)](#)]

# Seismic Newtonian noise

# How is seismic NN different from “seismic noise”?



# Ground model

Ground is modeled as a **linear elastic medium**; seismic waves are weak perturbations to the steady-state

The simplest possible ground model is a **homogeneous, isotropic, infinite half-space**, which is mechanically characterized by only three numbers:

- Bulk modulus,  $K$ : stiffness in response to pressure — typically ~1 MPa (soils) to >10 GPa (rocks)
- Shear modulus,  $\mu$ : stiffness in response to shear — typically ~1 MPa (soils) to >10 GPa (rocks)
- Density,  $\rho_0$ : typically several thousand kg/m<sup>3</sup>



# Seismic waves

Body waves: travel through the bulk of the ground

- Pressure (P) waves
- Shear (S) waves — two polarizations

Surface waves:

- Rayleigh waves
- (Love waves — do not appear in homogeneous media, and don't make NN in a half-space model)

Waves described by a displacement field  $\xi(\mathbf{r}, t)$ , which is connected to density fluctuations via  $\delta\rho(\mathbf{r}, t) = -\nabla \cdot [\rho_0(\mathbf{r}, t)\xi(\mathbf{r}, t)]$

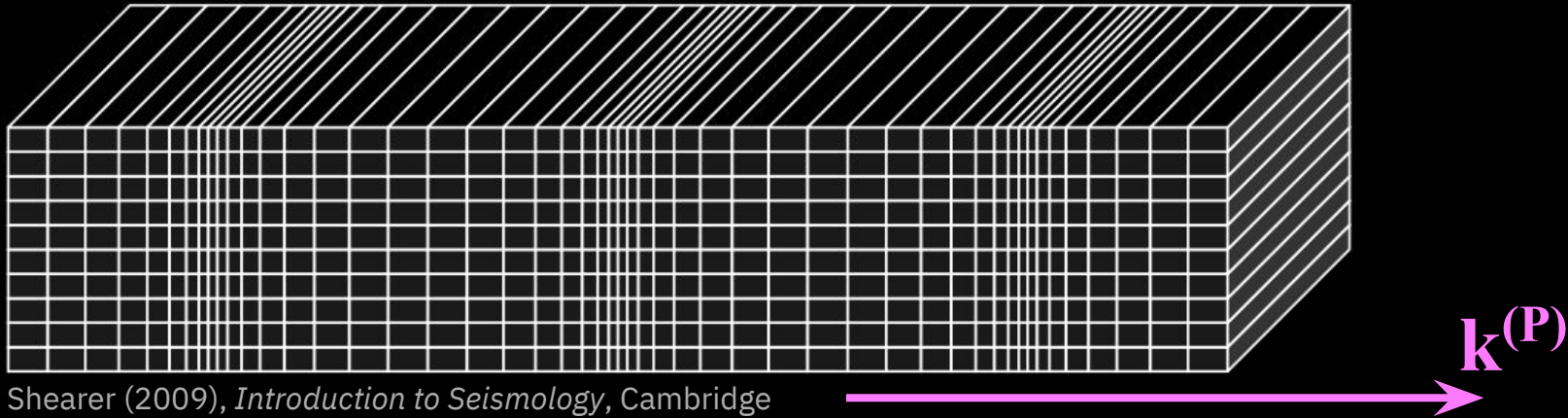
# Seismic P-waves

Longitudinal displacement field:  $\xi^{(P)}(\mathbf{r}, t) = \xi_0^{(P)} \cos(\mathbf{k}^{(P)} \cdot \mathbf{r} - \omega t) \hat{\mathbf{k}}$

Creates NN primarily through density perturbation

Wave speed:

$$c_P = \sqrt{\frac{K + 4\mu/3}{\rho_0}}$$



# Seismic S-waves

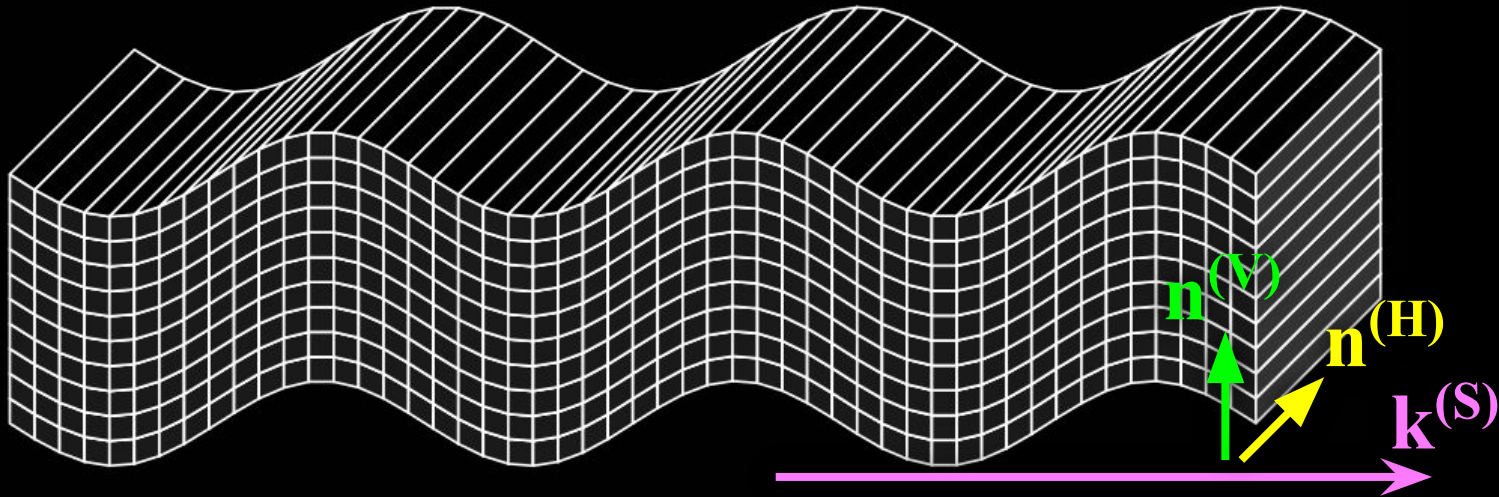
Transverse displacement field, two polarizations:

$$\xi^{(\text{SH})}(\mathbf{r}, t) = \xi_0^{(\text{SH})} \cos(\mathbf{k}^{(\text{S})} \cdot \mathbf{r} - \omega t) \hat{\mathbf{n}}^{(\text{H})}$$

$$\xi^{(\text{SV})}(\mathbf{r}, t) = \xi_0^{(\text{SV})} \cos(\mathbf{k}^{(\text{S})} \cdot \mathbf{r} - \omega t) \hat{\mathbf{n}}^{(\text{V})}$$

Motion is pure shear (no density changes). S-waves produce NN only by displacing boundaries.

$$\text{Wave speed } c_S = \sqrt{\mu/\rho_0}$$

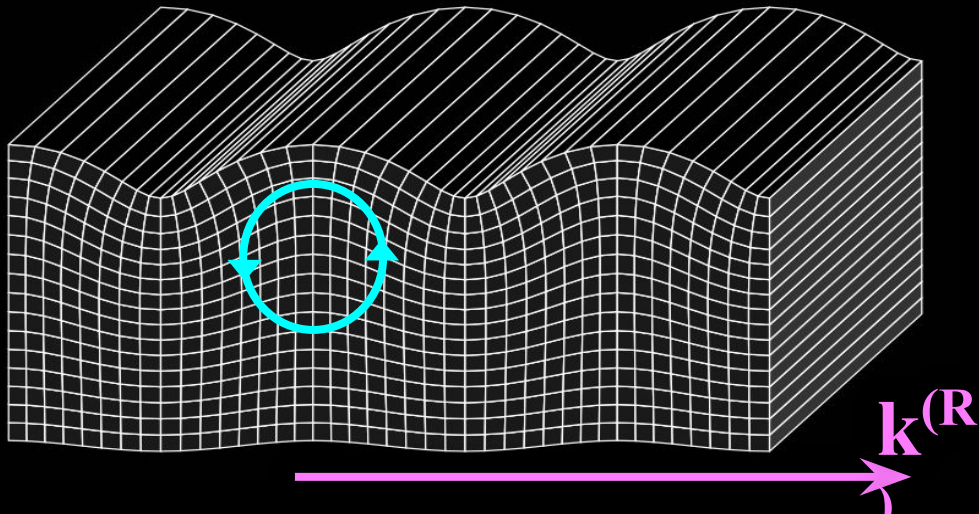


# Seismic Rayleigh waves

Surface acoustic wave; displacement field decays exponentially with depth

Retrograde elliptical particle motion: **creates NN through both density perturbation and ground surface displacement**

Travel at ~90% the speed of S-waves.



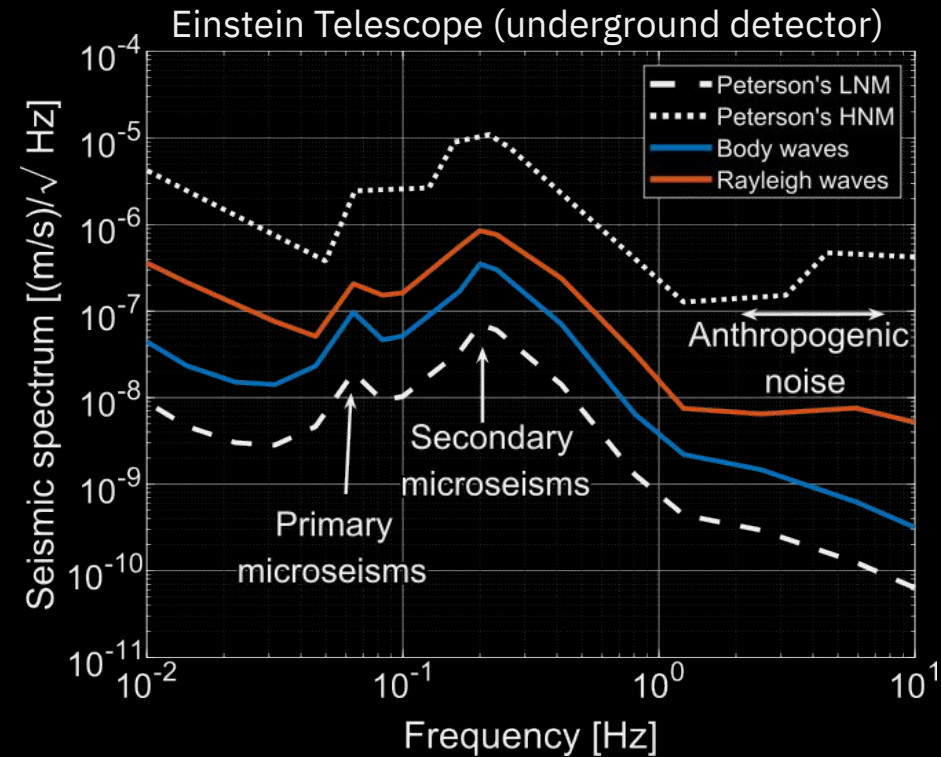
# Estimating seismic Newtonian noise

1. Establish a model for the ground.
2. Write down the displacement field  $\xi(\mathbf{r}, t)$  of each seismic wave.
3. Compute the corresponding density fluctuation and boundary perturbation.
4. Compute the local gravity fluctuation  $\delta\mathbf{g}(t)$  due to each wave, according to Newton's law of universal gravitation.
5. Account for the full seismic field by summing over the superposition of many individual seismic waves.

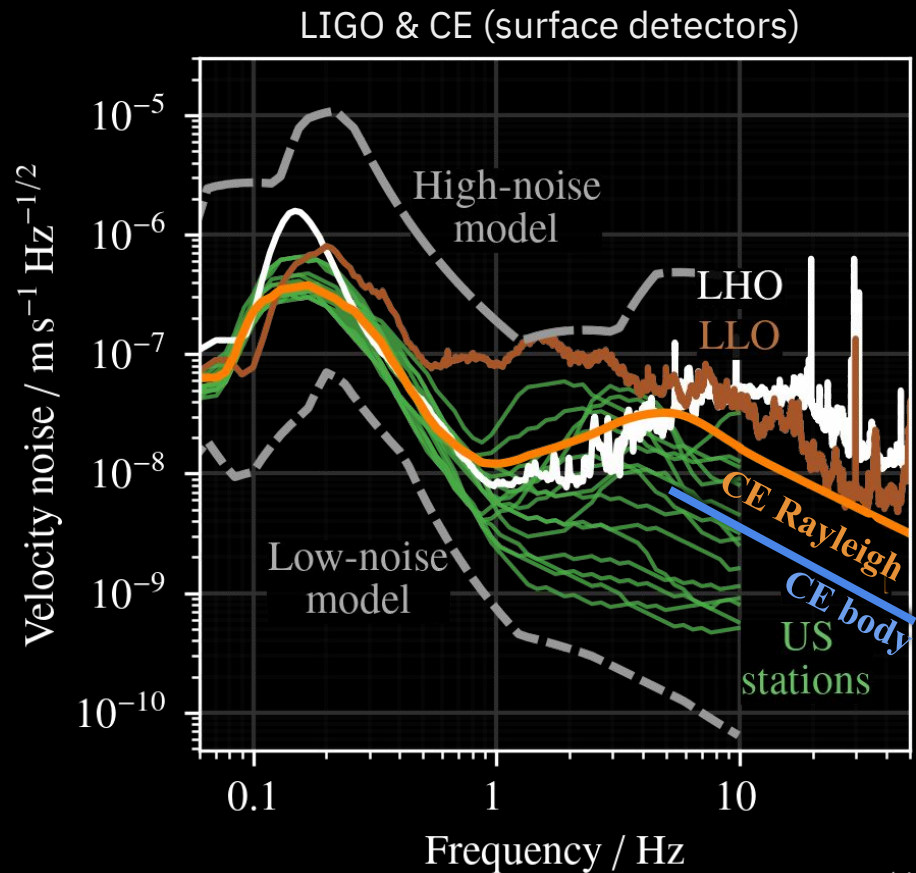
We typically start with an assumption of a **homogeneous, isotropic** seismic field.

In any case, we need to assume something about the body- and surface-wave content of the field...

# Seismic motion



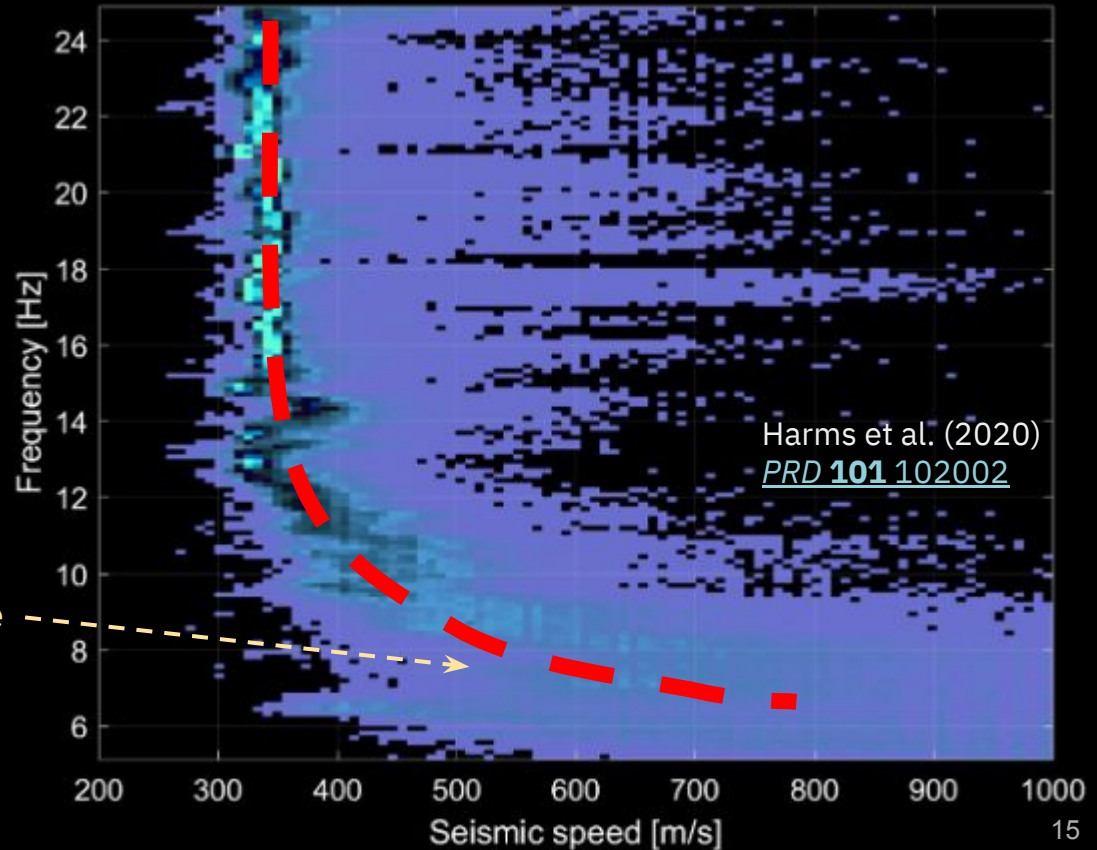
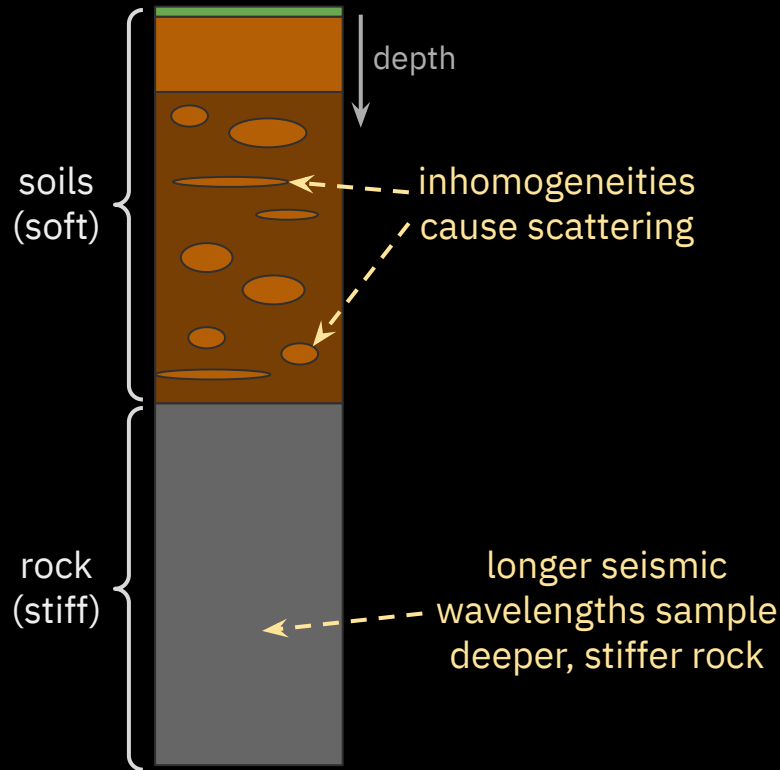
Amann et al. (2020) *RSI* **91**, 094504



Hall et al. (2021) arXiv:2012.03608

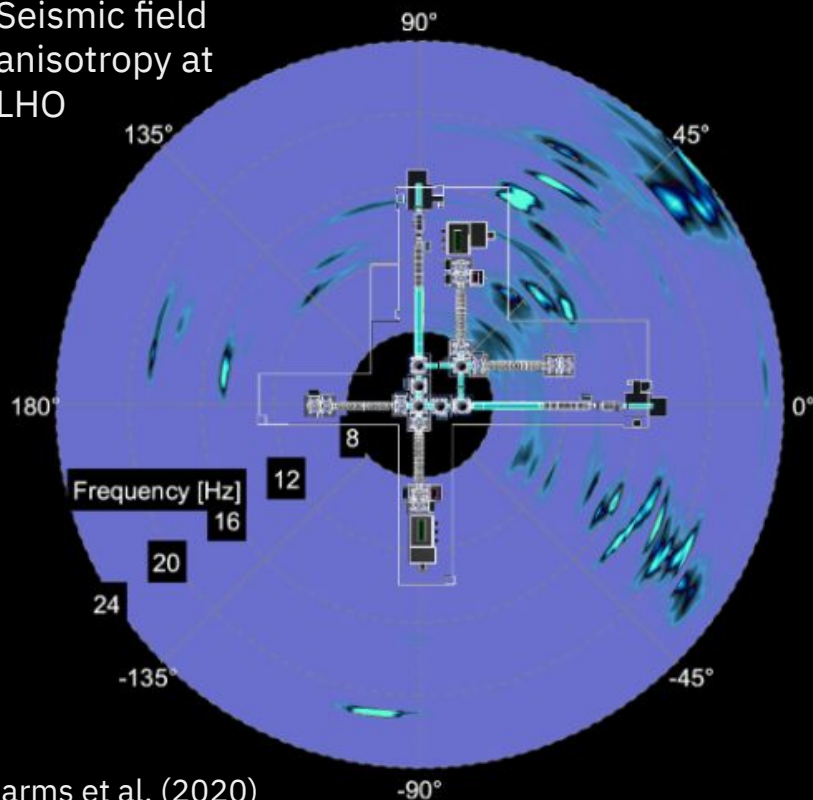
# Beyond the homogenous, isotropic ground model...

Rayleigh-wave dispersion at LHO



# Beyond the homogeneous, isotropic seismic field...

Seismic field  
anisotropy at  
LHO



Seismic fields (at least on the surface) are not isotropic.

In current detectors, most of the seismic noise is generated by the GW facility itself, sometimes very close to the test masses.



# Mitigation: site selection and detector geometry

0. Go into space

1. Pick a low-noise site

Can be 100× variation in typical seismic amplitude at 10 Hz between sites

2. Keep the site low noise

Isolate vibrating machinery from the ground near the test masses

3. Make the detector longer

Equivalent strain amplitude from NN scales like  $1/L$ .

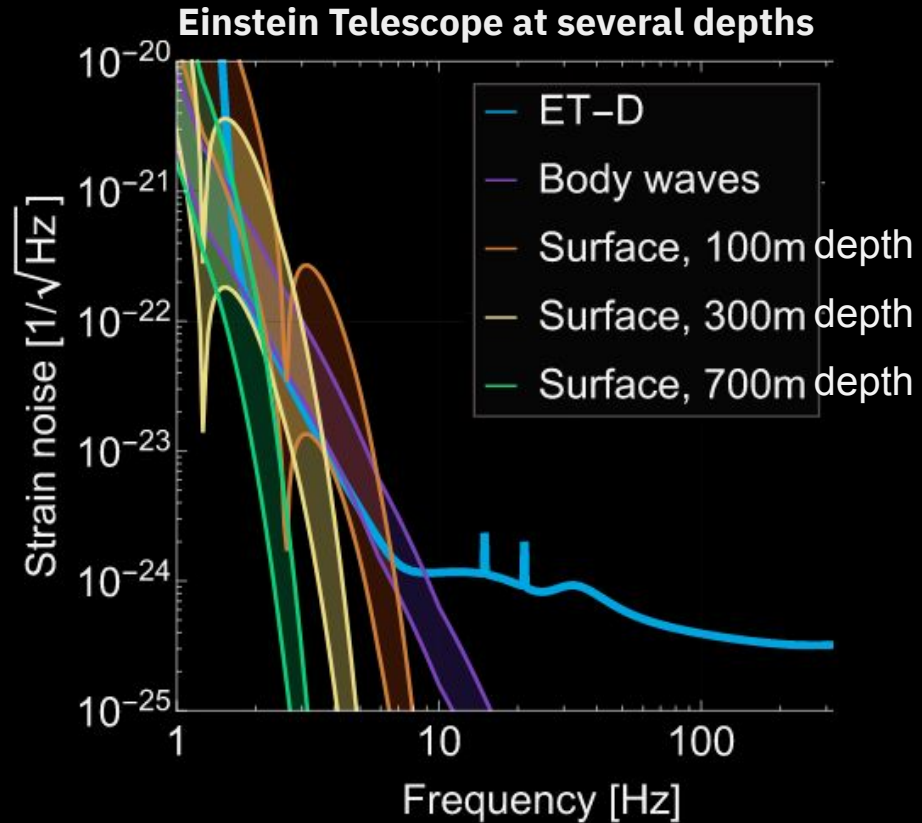
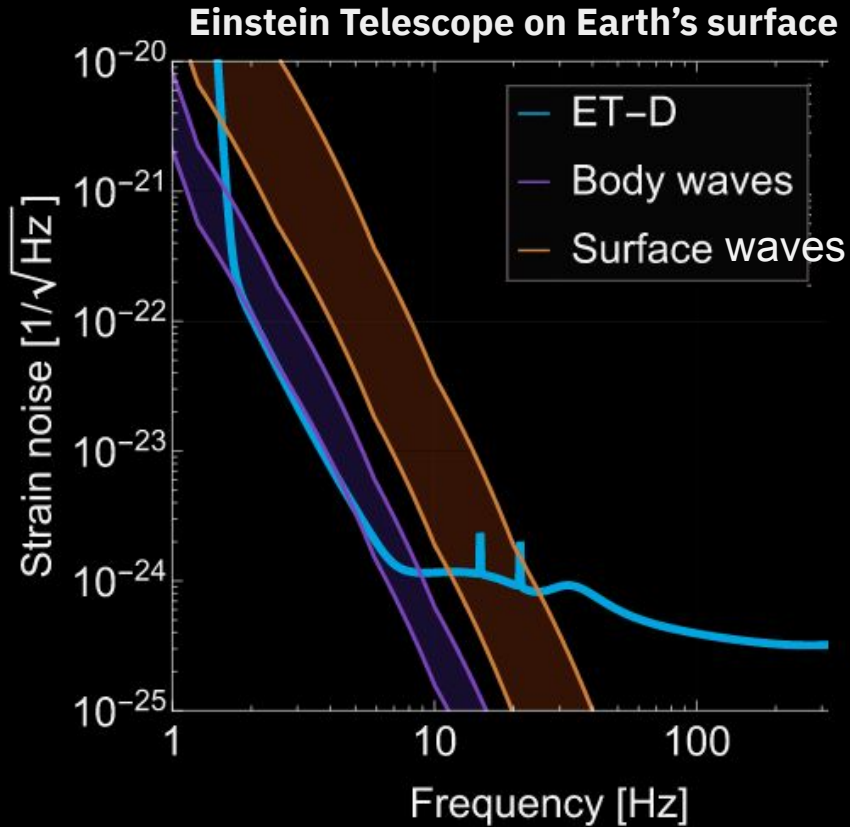
4. Go underground

Surface motion dominated by Rayleigh waves, which decay exponentially with depth.

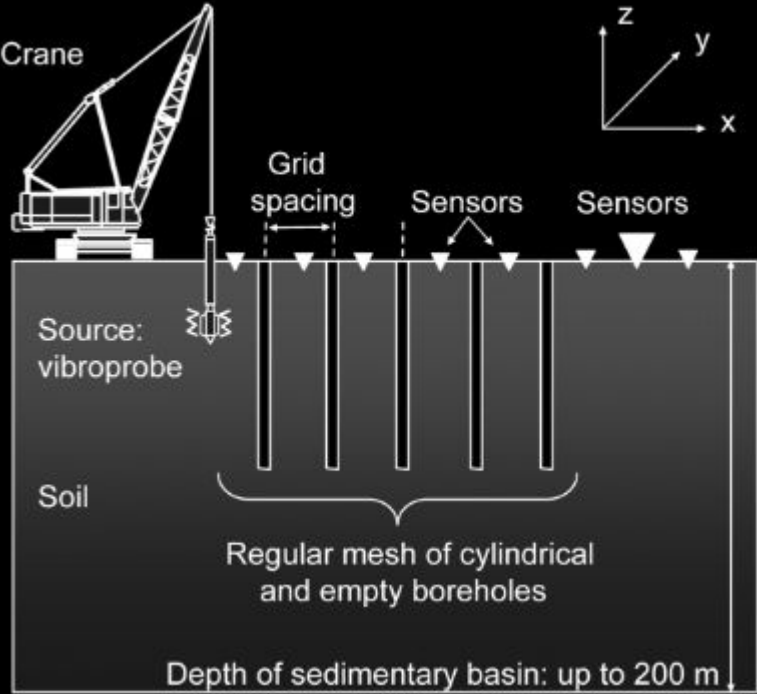
5. If on the surface, place test masses further away from the ground

Ditches, trenches, basements, geofoam...

# Mitigation: going underground



# Mitigation: seismic metamaterials

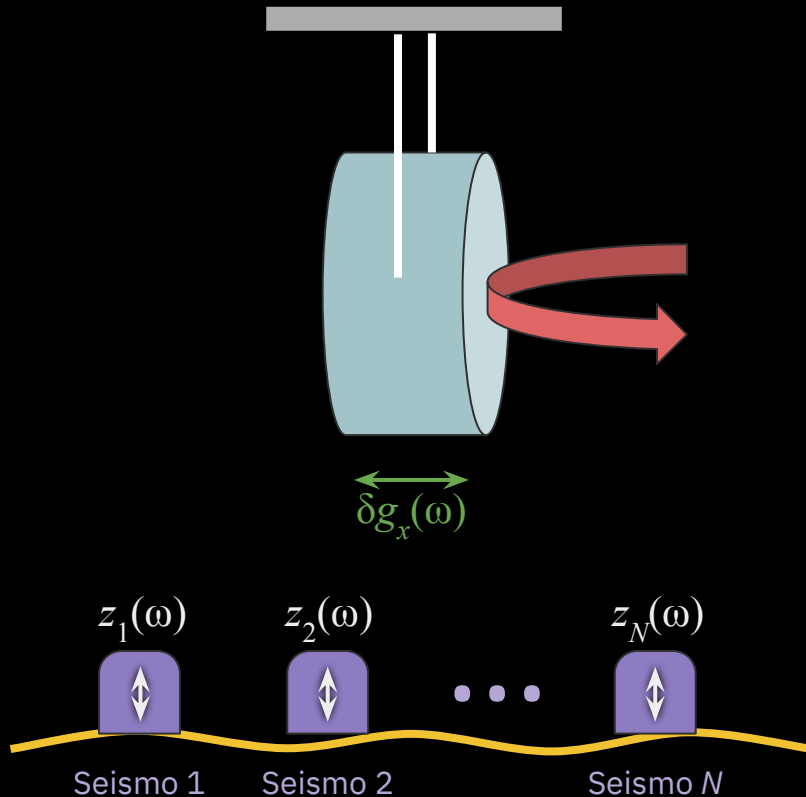


Sensitive three components velocimeters (green grid)

Five meters deep 320 mm holes

Source :  
 - Frequency : 50 Hz  
 - Horizontal displacement : 14 mm

# Mitigation: seismometer array subtraction



**Goal:** use a network of seismometers to estimate the local gravity fluctuation at the test mass.

Often formulated as a **Wiener filter problem**: given  $N$  seismometer data records  $\{z_i(\omega)\}$ , find a set of filters  $\{W_i(\omega)\}$  such that the sum  $\sum W_i(\omega)z_i(\omega)$  is an optimal estimate of the local gravity  $\delta g_x(\omega)$  at the test mass.

“Optimal” here means minimizing the mean-squared error between the estimate and the true value of  $\delta g_x(\omega)$ .

## Mitigation: seismic array subtraction (2)

Wiener filter solution requires measuring (or modeling):

1. the  $N$  cross-spectra  $\{H_i(\omega)\}$  from each seismometer to the test mass local gravity,
2. and the  $N \times N$  cross-spectra  $\{C_{ij}(\omega)\}$  between all the seismometers.

Then the optimal filters are  $W_i = H_j(C^{-1})_{ij}$ , and the residual mean-squared error is

$$R(\omega) = 1 - \frac{\mathbf{H}(\omega)^T \mathbf{C}(\omega)^{-1} \mathbf{H}(\omega)}{S_{\delta g}(\omega)}$$

where  $S_{\delta g}(\omega)$  is the power spectrum of the Newtonian noise in the detector.

# Mitigation: seismic array subtraction (3)

**Noise limit:**  $R(\omega) > 1/(N \cdot \text{SNR}(\omega)^2)$

- Typical seismometer SNR ~ 100 to 1000
- Typical array size of order 10

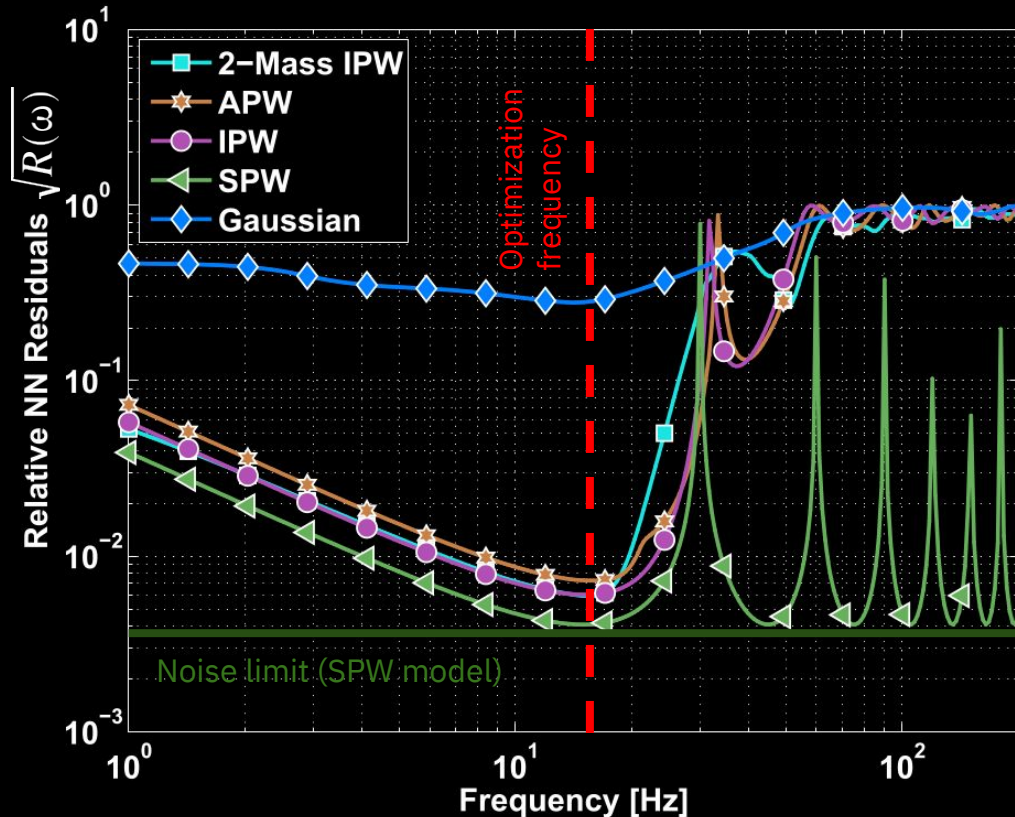
**Geometric limit:** depends on the seismometer coordinates and the seismic field

- Array scale should match the seismic wavelength of interest
- Array coordinates must be optimized, generally numerically

Driggers et al. (2012) [PRD 86 102001](#)

$$R(\omega) = 1 - \frac{\mathbf{H}(\omega)^T \mathbf{C}(\omega)^{-1} \mathbf{H}(\omega)}{S_{\delta g}(\omega)}$$

# Example: aLIGO Rayleigh-wave NN cancellation model



In simulation, 100× cancellation of Rayleigh waves seems feasible on the surface, over a limited frequency band.

Also in simulation, 10× cancellation of body waves seems feasible underground.

[Badaracco & Harms (2019), [COG 36 145006](#)]

(A more cautious assumption might be 10× Rayleigh wave cancellation and 3× body wave cancellation...)

# Seismic NN: what is computed in pygwinc?

`newtonian.gravg_rayleigh`

`newtonian.gravg_pwave`

`newtonian.gravg_swave`

Assumptions:

1. Ground is a homogenous, isotropic, infinite half-space
2. Seismic field is a stationary, homogeneous, isotropic superposition of plane waves arriving from all directions ( $2\pi$  sr for P- and S-waves;  $2\pi$  rad for Rayleigh waves)
3. For Rayleigh waves, the ground motion is an ad-hoc broken power law with an optional constant describing the amount of subtraction
4. For body waves, the ground motion is specified as a multiple of the Peterson low-noise model
5. Scattering/interconversion of seismic waves is ignored



# Example (from Cosmic Explorer 2)

In the `ifo.yaml` file...

```
66 Seismic:
67   Site: 'LHO' # LHO or LLO (only used for Newtonian noise)
68   KneeFrequency: 5 # Hz; freq where 'flat' noise rolls off
69   LowFrequencyLevel: 1e-9 # m/rHz; seismic noise level below f_knee
70   KneeFrequencyHorizontal: 4 # Hz; freq where 'flat' noise rolls off
71   LowFrequencyLevelHorizontal: 1e-9 # m/rHz; seismic noise level below f_knee
72   Gamma: 0.8 # abruptness of change at f_knee
73   Rho: 1.8e3 # kg/m^3; density of the ground nearby
74   Beta: 0.8 # quiet times beta: 0.35-0.60
75             # noisy times beta: 0.15-1.4
76   Omicron: 10 # Feedforward cancellation factor
77   TestMassHeight: 1.5 # m
78   pWaveSpeed: 600 # m/s
79   sWaveSpeed: 300 # m/s
80   RayleighWaveSpeed: 250 # m/s
81   pWaveLevel: 15 # Multiple of the Peterson NLNM amplitude
82   sWaveLevel: 15 # Multiple of the Peterson NLNM amplitude
83   PlatformMotion: '6D'
```

ground motion model  
used by `gravg_rayleigh`

ground motion model  
used by `gravg_pwave` and  
`gravg_swave`

# Atmospheric Newtonian noise

# Infrasound NN

Sound (acoustic waves) below 20 Hz

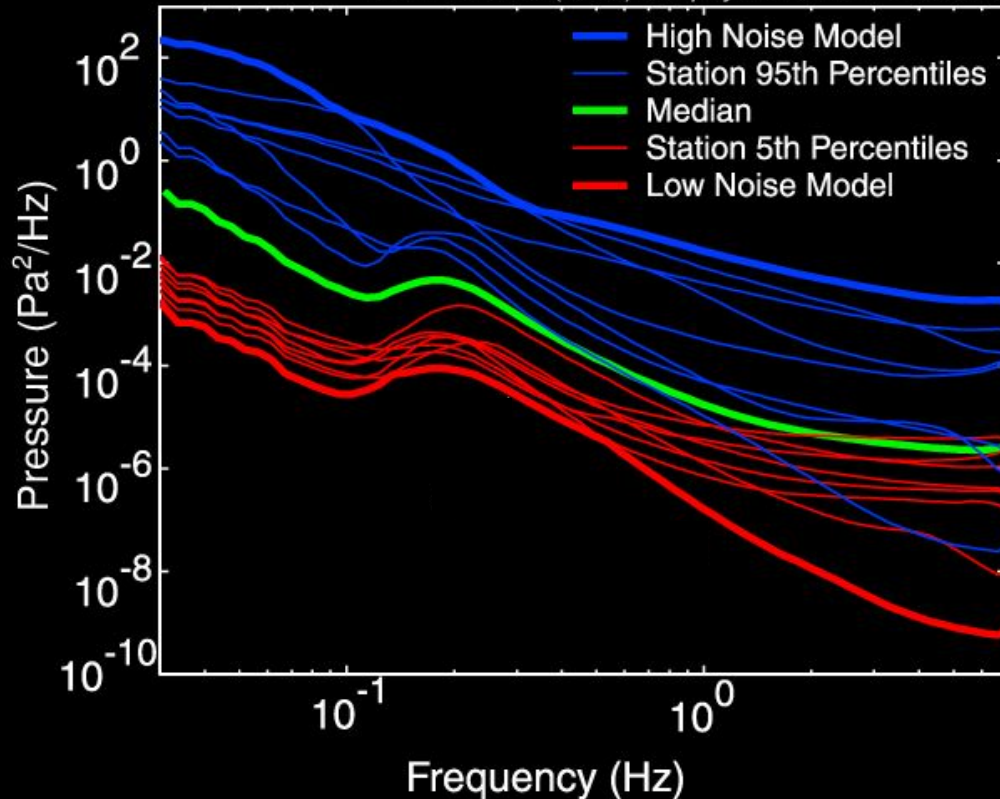
Pressure field  $p(\mathbf{r}, t) = p_0 \cos(\mathbf{k} \cdot \mathbf{r} - \omega t)$ ; speed  $c_{\text{sound}} = 340 \text{ m/s}$

Produces density fluctuation via  $\gamma(\delta\rho/\rho) = (\delta p/p)$ ;  $\gamma = 1.4$  for air

NN calculations often assume a homogeneous, isotropic acoustic wave field, but real wavefields can be anisotropic.

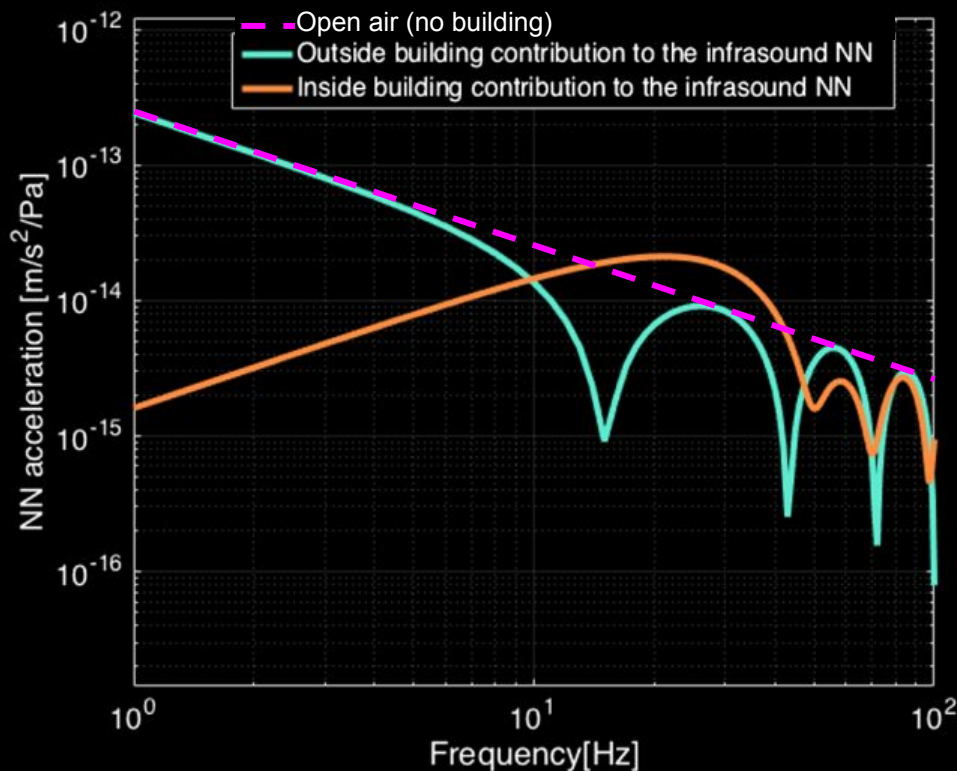
# Infrasound spectrum

Bowman et al. (2005) *Geophys Res Lett* **32** L09803



Median model suggests  $1 \text{ mPa}/\text{Hz}^{1/2}$  above a few hertz, but potentially limited by sensor noise or calibration.

# Structural mitigation of infrasonic NN?



An exponential suppression of outside infrasound NN requires going underground: at a depth  $d$ , the cutoff frequency is  $f = c_{\text{sound}}/d$

Buildings and burial do not save you from noise generated from indoor detector infrastructure, like HVAC systems.

# Subtracting infrasonic NN?

Navier–Stokes equations are highly nonlinear, leading to contributions to the pressure field from **turbulent flow**, particularly during windy times:

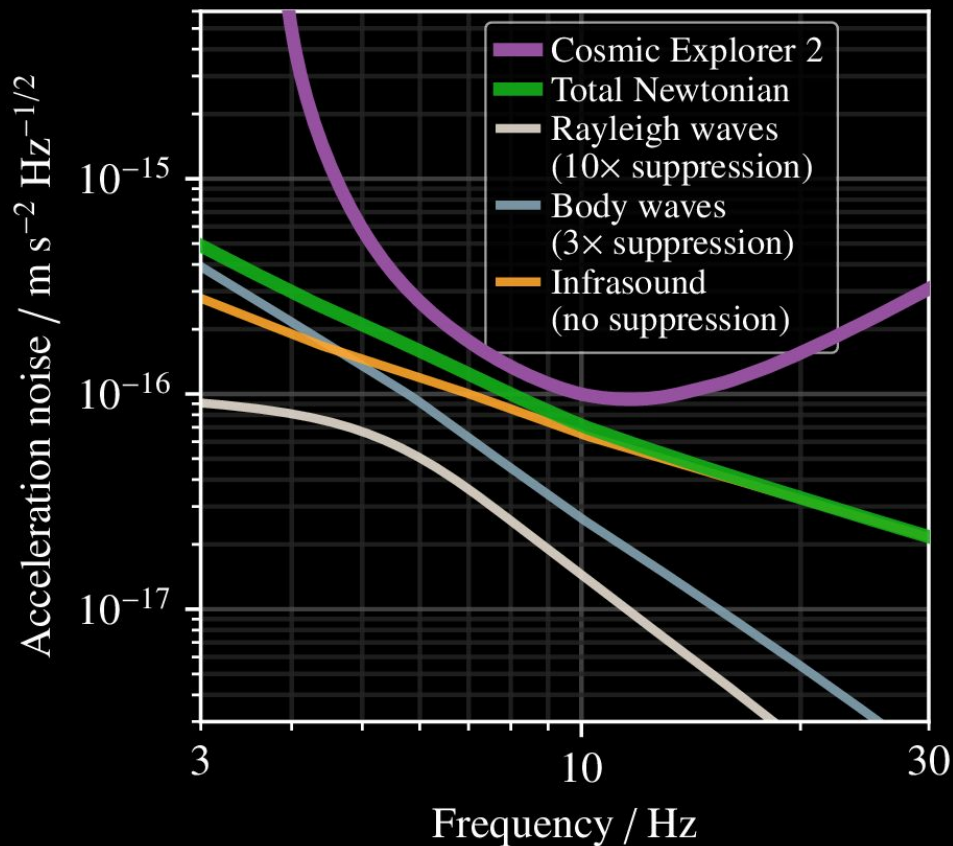
- Intrinsic wind noise (turbulence–turbulence and turbulence–shear interactions)
- Stagnation pressure: wind blowing against the sensor.

Raspet et al. (2019) in *Infrasound Monitoring for Atmospheric Studies* (Le Pichon et al., eds.), Springer.

LIDAR? Currently can sense  $\delta p/p \sim 10^{-3}$ , but infrasound NN fluctuations are  $\sim 10^{-7}$ .

Fiorucci et al. (2018) [PRD 97 062003](#)

# NN in a future surface detector



Assumed  $1 \text{ mPa/Hz}^{1/2}$  pressure spectrum; no consideration of building.

# Other atmospheric NN sources

Inhomogeneities of air temperature and humidity, which are advected past the test mass by wind:

- Exponentially suppressed above a few hertz for typical surface detector parameters (wind speeds  $\sim 10$  m/s, building size  $\sim 10$  m)

Aeroacoustic noise (Lighthill noise):

- Pressure fluctuations sourced by turbulent flow
- Also unlikely to be significant above a few hertz



# What is computed in pygwinc?

`newtonian.atmois` (infrasonic NN)

Assumptions:

- Stationary, homogeneous, isotropic superposition of plane waves
- Pressure field can be specified as a power law, or else as the global median (Bowman) model
- Ground is perfectly reflecting (no seismic/infrasound interconversion)

# Example (from Cosmic Explorer 2)

In the `ifo.yaml` file...

```
86 Atmospheric:
87   AirPressure: 101325          # Pa
88   AirDensity: 1.225           # kg/m**3
89   AirKinematicViscosity: 1.8e-5 # m**2/s
90   AdiabaticIndex: 1.4        #
91   SoundSpeed: 344            # m/s
92   WindSpeed: 5                # m/s; typical value
93   Temperature: 300           # K
94   TempStructConst: 0.2        # K**2/m**(2/3);
95   TempStructExp: 0.667        #
96   TurbOuterScale: 100        # m
97   # TurbEnergyDissRate: 0.01  # m**2/s**3
98   KolmEnergylm: 1            # Kolmogorov energy spectrum at 1/m [m**2/s**2]
```

# Homework

Jan Harms (2019), “Terrestrial gravity fluctuations”, *Living Reviews in Relativity* **22**(6), <https://doi.org/10.1007/s41114-019-0022-2>

Read through 3.2.1.:

- Understand how to formulate Newton’s law in terms of gravitational potential, and how to separate potential fluctuations into to bulk and surface contributions.

Read through 3.4.2.:

- Familiarize yourself with the displacement fields for Rayleigh waves (Eq. 91, referring back to Eqs. 37–41 and related discussion as necessary).
- Work through the subsequent calculation for the gravitational potential fluctuations for a Rayleigh wave.