## **Gravitational Wave Transients**

Ian Jones

Build-a-detector workshop

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# What is a "gravitational wave transient"

- No universally accepted definition.
- Let's go for:

"Any gravitational wave signal of duration less than  $\sim$  1 day, excluding signals from the inspiral/merger/ring-down of compact binary coalescence"

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# Examples LSC/LVC/LVK papers on transients

#### Of $\sim$ 210 published LSC/LVC/LVK papers, approx 47 target transients. Examples:

- Optically targeted search for gravitational waves emitted by core-collapse supernovae during the first and second observing runs of advanced LIGO and advanced Virgo" (2020).
- 2 "Search for Gravitational-wave Signals Associated with Gamma-Ray Bursts during the Second Observing Run of Advanced LIGO and Advanced Virgo" (2019).
  - 3) "All-sky search for short gravitational-wave bursts in the second Advanced LIGO and Advanced Virgo run" (2019).
  - 4) "All-sky search for long-duration gravitational-wave transients in the second Advanced LIGO observing run" (2019).
- Search for Gravitational Waves from a Long-lived Remnant of the Binary Neutron Star Merger GW170817" (2019).
- Search for Transient Gravitational-wave Signals Associated with Magnetar Bursts during Advanced LIGO's Second Observing Run" (2019).
- Constraints on cosmic strings using data from the first Advanced LIGO observing run" (2018).
- Search for transient gravitational waves in coincidence with short-duration radio transients during 2007-2013" (2016).
- "Multimessenger search for sources of gravitational waves and high-energy neutrinos: Initial results for LIGO-Virgo and IceCube" (2014).
  - "Implications for the Origin of GRB 051103 from LIGO Observations" (2012).
  - "Search for gravitational waves associated with the August 2006 timing glitch of the Vela pulsar" (2011).

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# Types of transient

- Let's divide into three classes:
  - Catastrophic collapse
  - Cosmic strings
  - Oscillating/rotating neutron stars

We'll mention all three, but focus on the third.

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## Stellar collapse

- Supernova explosion mechanism not well understood, but collapse will be asymmetric at some level.
- Most recent LVK paper targeted 5 optically observed SN, all quite close (7–20 Mpc).
- Predicted peak frequency for GW emission uncertain:

| Waveform family                                      | Waveform identifier      | $h_{\rm rss}~(10^{-22}\sqrt{\rm s}@10~{\rm kpc})$ | $f_{\text{peak}}$ (Hz) | $E_{\rm GW}~(10^{-9}~M_{\odot}c^2)$ | Polarizations |
|--|--------------------------|---|------------------------|-------------------------------------|---------------|
| Müller et al. [95] 3D<br>convection and SASI         | mul1-L15-3               | 1.655   | 150                    | $3.741 \times 10^{-2}$              | +, ×          |
|  | mul2-N20-2               | 3.852   | 176                    | $4.370 \times 10^{-2}$              | +, ×          |
|  | mul3-W15-4               | 1.093   | 204                    | $3.247\times 10^{-2}$               | +, ×          |
| Ott et al. [96] 3D<br>convection and SASI            | ott1-s27fheat1p05        | 0.238   | 1019                   | $7.342\times10^{-1}$                | +, $\times$   |
| Yakunin <i>et al.</i> [97] 2D<br>convection and SASI | yak1-B12-WH07            | 3.092   | 760                    | 3.411                               | +             |
|  | yak2-B15-WH07            | 14.16   | 932                    | 7.966                               | +             |
|  | yak3-B20-WH07            | 3.244   | 638                    | 4.185                               | +             |
|  | yak4-B25-WH07            | 18.05   | 1030                   | 14.21                               | +             |
| Scheidegger et al. [98]<br>rotating core collapse    | sch1-R1E1CAL             | 0.129   | 1155                   | $1.509 \times 10^{-1}$              | +, ×          |
|  | sch2-R3E1AC <sub>1</sub> | 5.144   | 466                    | $2.249 \times 10^{2}$               | +, ×          |
|  | sch3-R4E1FCL             | 5.796   | 698                    | $4.023 \times 10^2$                 | $+, \times$   |
| Dimmelmeier et al. [99]<br>rotating core collapse    | dim1-s15A2O05ls          | 1.052   | 770                    | 7.685                               | +             |
|  | dim2-s15A2O09ls          | 1.803   | 754                    | 27.880                              | +             |
|  | dim3-s15A3O15ls          | 2.690   | 237                    | 1.380                               | +             |

Table from Abbott+ PRD 101 084002 (2020)

## Stellar collapse

Energetics suggest that we'd need an extremely energetic and nearby event to make a detection:



Table from Abbott+ PRD 101 084002 (2020)

## **Cosmic strings**

- Can have GW emission from cusps/kinks.
- Large range of GW frequencies relevant; most recent search covered full LIGO/Virgo band.



Table from Abbott+ PRD 97 102002 (2018)

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# Neutron stars and oscillations

- Young pulsars glitch.
- Magnetars flare.
- Low-mass X-ray binaries burst.





Crab pulsar

SGR 1900+14

LMXB

• Which sorts of oscillation mode are relevant, to what amplitudes are they excited, and what are the GW/EM signitures?

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# Types of neutron star oscillation modes

Neutron stars can oscillate, emitting GWs in the process.

- There exist a whole zoo of types of these normal modes:
  - f-mode: the fundamental mode (restoring force is pressure).
  - p-mode: overtones of fundamental (restoring force is pressure).
  - g-mode: restoring force due to composition/temperature gradients.
  - r-mode: restoring force in Coriolis force (only in rotating stars).
  - Alfvén mode: restoring force due to magnetic field.
  - Elastic mode: restoring force elastic.
- Key point is that frequency and lifetime of oscillation related to internal stellar structure.
- Hope to perform 'gravitational wave asteroseismology'!

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## **Basic equations**

- Should work in General Relativity, but here's a Newtonian treatment, for simplicity.
- Fundamental equations relate pressure P, velocity v, density ρ and gravitational potential
   φ:
  - Newton's 2nd law:

$$\frac{d\mathbf{v}}{dt} = -\frac{1}{\rho}\nabla P - \nabla \Phi + \mathbf{F}_{\text{other}},\tag{1}$$

Poisson's equation:

$$\nabla^2 \Phi = 4\pi G\rho, \tag{2}$$

Conservation of mass:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = \mathbf{0},\tag{3}$$

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where  $\mathbf{F}_{\text{other}}$  contains any other forces you choose to model (viscosity, magnetic fields, elasticity, . . . )

## Computing modes

If unperturbed star is static and spherical, we can write

$$P_{\text{perturbed}}(t, r, \theta, \phi) = P_{\text{unperturbed}}(r) + \delta P(t, r, \theta, \phi).$$
(4)

- Can then substitute into the differential equations, and separate out the zero order (unperturbed) terms and the first order terms.
- Typically expand perturbed quantities assuming a harmonic time dependence, and in terms of spherical harmonics, e.g.

$$\delta P(t, r, \theta, \phi) = \sum_{lm} \delta P_{lm}(r) Y_{lm}(\theta, \phi) e^{i\omega_{lm}t}.$$
(5)

 Imposing the boundary conditions then leads to eigenmode solutions with discrete frequencies.

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## Simplest example: non-rotating uniform density star

• For non-rotating uniform density fluid star, can analytically find 'Kelvin modes', with frequency

$$\omega_{lm}^2 = \frac{4\pi G\rho}{3} \frac{2l(l-1)}{2l+1}.$$
 (6)

• Writing in terms of  $f_{lm} = \omega_{lm}/(2\pi)$  and parameterising:

$$f_{lm} \approx 3 \,\mathrm{kHz} \, \sqrt{\frac{l(l-1)}{2l+1}} \left(\frac{M}{1.4M_{\odot}}\right)^{1/2} \left(\frac{10^6 \,\mathrm{cm}}{R}\right)^{3/2}.$$
 (7)

e.g. for *I* = 2, get *f* ≈ 2 kHz. This sets the frequency scale for GW for such modes—at the extreme high end of LIGO/Virgo.

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#### Glitches

- Most of the time pulsar spin frequencies gradually decrease.
- Occasionally some younger pulsars undergo sudden spin-ups.



Fractional variation in spin frequency small, e.g.

$$\frac{\Delta\Omega}{\Omega}\sim 10^{-6}$$

for Vela; this is considered a violent glitchier!

Key question: do glitches generate GWs?

### Glitch energies: a naive estimate

In absence of detailed model, can make a 'naive' estimate:

$$E_{\text{glitch}} = I\Omega\Delta\Omega = I\Omega^2 \frac{\Delta\Omega}{\Omega}.$$

Parameterising with Vela in mind:

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$$\begin{split} E_{\rm glitch} \sim 4.95 \times 10^{42} \, \mathrm{erg} \, \left(\frac{f_{\rm spin}}{11.2 \, \mathrm{Hz}}\right)^2 \left(\frac{\Delta \Omega / \Omega}{10^{-6}}\right), \\ \eta_{\rm rss} \sim 10^{-22} \, \mathrm{Hz}^{1/2} \left(\frac{287 \, \mathrm{pc}}{r}\right) \left(\frac{f_{\rm spin}}{11.2 \mathrm{Hz}}\right) \left(\frac{1 \, \mathrm{kHz}}{f_{\rm GW}}\right) \left(\frac{\Delta \Omega / \Omega}{10^{-6}}\right)^{1/2} \end{split}$$

But can all this energy be put into modes, and if so which?

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## The Vela glitch paper: f-modes

- Search for 'fundamental' f-mode excitation following a Vela glitch was carried out (Abadie+ 2011).
- Looked for damped sinusoids with  $f \sim 1-3$  kHz,  $\tau \lesssim 0.5$  s.
- Found energy release  $\Delta E_{\rm GW} \lesssim 10^{45}$  erg.



# Do glitches excite oscillation modes?

There have been a small number of attempts to answer this question.

- van Eysden and Melatos (2008) looked at two-component infinitely long cylinder.
- Sidery, Passamonti, Andersson (2009) looked at two-component spherical spin-up using time evolutions.
- Keer & DIJ (2015) looked at the *starquake* model.

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# Starquakes & mode excitaiton

Keer & DIJ (2015) considered a specific evolutionary sequence:



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# Projecting the initial data onto the modes

Do indeed find that fundamental-like mode excited:



## Magnetars

- Magnetars are neutron stars with very strong magnetic fields,  $B \sim 10^{15}$  G.
- Two sub-classes seem to occasionally emit violent outbursts of gamma rays:
  - Soft-gamma ray pulsars (SGRs)
  - Anomalous X-ray pulsars (AXPs)
- Believed to represent reconfigurations of the magnetic field, possibly triggered by crust cracking.
- Oscillations seen in the light curves of some bursts:



Figure credit Watts & Strohmayer (2006)

#### Magnetars

- Note clear what the excited modes are.
- Frequencies from tens of Hz to hundreds.
- Could be some combination of elastic and magnetic forces.
- Example: LVK searched over wide range for February 25 Burst from SGR 1806-20:

| Frequency (Hz) | Tau (s) | h <sub>rss</sub>       |                        | Energy (erg)          |                       |
|----------------|---------|------------------------|------------------------|-----------------------|-----------------------|
|                |         | Half Sine-Gaussian     | Ringdown               | Half Sine-Gaussian    | Ringdown              |
| 55             | 400     | $2.29 \times 10^{-22}$ | $2.43 \times 10^{-22}$ | $1.82 \times 10^{44}$ | $2.06 \times 10^{44}$ |
| 55             | 150     | $1.97 \times 10^{-22}$ | $2.11 \times 10^{-22}$ | $1.35 \times 10^{44}$ | $1.55 \times 10^{44}$ |
| 150            | 400     | $1.32 \times 10^{-22}$ | $1.37 \times 10^{-22}$ | $4.52 \times 10^{44}$ | $4.86 \times 10^{44}$ |
| 150            | 150     | $1.14 \times 10^{-22}$ | $1.22 \times 10^{-22}$ | $3.37 \times 10^{44}$ | $3.89 \times 10^{44}$ |
| 450            | 400     | $1.69 \times 10^{-22}$ | $1.79 \times 10^{-22}$ | $6.62 \times 10^{45}$ | $7.47 \times 10^{45}$ |
| 450            | 150     | $1.78 \times 10^{-22}$ | $1.83 \times 10^{-22}$ | $7.43 \times 10^{45}$ | $7.83 \times 10^{45}$ |
| 750            | 400     | $2.56 \times 10^{-22}$ | $2.70 \times 10^{-22}$ | $4.21 \times 10^{46}$ | $4.69 \times 10^{46}$ |
| 750            | 150     | $2.11 \times 10^{-22}$ | $2.37 \times 10^{-22}$ | $2.87 \times 10^{46}$ | $3.61 \times 10^{46}$ |
| 1550           | 400     | $5.86 \times 10^{-22}$ | $6.22 \times 10^{-22}$ | $9.21 \times 10^{47}$ | $1.03 \times 10^{48}$ |
| 1550           | 150     | $4.38 \times 10^{-22}$ | $4.58 \times 10^{-22}$ | $5.16 \times 10^{47}$ | $5.62 \times 10^{47}$ |

Figure credit Abbott+ ApJ 874 63 (2019)

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## **R**-modes

- R-modes occur only in rotating stars; the restoring force is the Coriolis force.
- The GW frequency is proportional to the star's spin frequency:

$$f_{\rm GW} \approx \frac{4}{3} f_{\rm spin}.$$
 (8)

• The r-modes are prime candidates for undergoing the *Chandrasekhar-Friedman-Schutz* (CFS) instability.

R-mode image: Hannah/Owen:



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## Growth of the instability

• Amplitude of a normal mode evolves as  $\alpha(t) \sim e^{-t/\tau}$ , where



R-mode gravitational radiation reaction timescale scales sharply with f<sub>spin</sub>:

$$\tau_{\rm GRR} \approx -1 \text{ hour } \left(\frac{435 \,\text{Hz}}{f_{\rm spin}}\right)^6 \frac{1}{M_{1.4} R_{11.7}^4}$$

for an n = 1 polytrope (Owen et al 1998); weak dependence on equation of state.

Viscous damping timescale is temperature and spin frequency dependent:

$$\tau_{\rm damp} = \tau_{\rm damp}(f_{\rm spin}, T) > 0,$$

strongly dependent on details of the equation of state.

# Growth of the instability cont ...

For given temperature, mode goes 'CFS-unstable' for  $f_{spin} > f_{CFS}$ :



- Newly born neurtron star may be very hot with fast spin.
- Can then cooling into instability window, and spin down rapidly.
- There is a large body of literature debating the details of this!

# Summary

- There are many different sources for transient gravitational waves, including
  - Core collapse
  - Cosmic strings
  - Oscillating neurtron stars
- Frequency ranges of emission rather uncertain.
- For oscillations, frequency depends on sort of mode, and can be as high as several kHz.

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# Further reading

- Textbook: Shapiro & Teukolsky Black Holes, White Dwarfs, and Neutron Stars: The Physics of Compact Objects, especially chapters 6, 9, 10.
- Textbook: Andersson Gravitational-Wave Astronomy Exploring the Dark Side of the Universe, especially chapters 6, 13, 15, 18.
- Review article: Glampedakis & Gualtieri Gravitational waves from single neutron stars: an advanced detector era survey, arXiv 1709.07049.

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## Exercises

- Use dimensional analysis to estimate the dynamical timescale for a self-gravitating star of mass *M* and radius *R*. Parameterise your result in units of  $M = 1.4M_{\odot}$  and R = 10 km. This sets the frequency scale for f-mode oscillation, and also the Keplerian (i.e. break-up) frequency, which in turn sets the maximum r-mode frequency.
- ② Use dimensional analysis to estimate the timescale for oscillations in the elastic crust of a neutron star, with shear modulus  $\mu$ , mass *M* and radius *R*. Parameterise your answer in terms of  $\mu = 10^{30}$  erg cm<sup>-3</sup>,  $M = 1.4M_{\odot}$  and R = 10 km. This sets the frequency for torsional crust oscillations.
- **3** Use dimensional analysis to estimate the timescale for *Alfvén* (i.e. magnetic) oscillations for a star with a magnetic field of strength *B*, mass *M* and radius *R*. Parameterise in units of  $B = 10^{15}$  G,  $M = 1.4M_{\odot}$  and R = 10 km. [Hint: if you use Gaussian cgs units, then  $B^2$  (with *B* in units of Gauss) is an energy density, with units erg cm<sup>-3</sup>.